

Transport of perpendicular momentum across the last closed surface

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Questions

- What is the rate of perpendicular momentum transport across the last closed surface (LCS)?
 - characterize edge poloidal momentum “source”
- What is the role of Reynolds stress generation and transport by blobs?
 - Coppi [EPS 2006 paper P4.017 & NF 2002] blob momentum and recoil force, spontaneous toroidal rotation
- What determines the direction of the momentum transport (sign of edge source)?
- What is the role of sheath and volume dissipation?
- How does the dynamics of zonal momentum $\langle n_e v_y \rangle$ govern zonal flow $\langle v_y \rangle$ and turbulence regulation?
 - steep $n(x)$ at edge and $\delta n/n \sim 1$
- What can we learn from quasilinear averages?

Model

- minimal 2D turbulence model in plane normal to B for edge and SOL
 - Lodestar SOLT code
- equations:
 - continuity equation for density n
 - lost particles re-supplied at zero velocity
 - vorticity equation for “fluctuating” potential $\tilde{\Phi} = \Phi - \langle \Phi \rangle$
$$\langle \Phi \rangle = \int dy \Phi / L_y$$
 - curvature drive
 - adiabaticity parameter for electrons α_{dw} controls drift wave (DW) drive
 - sheath charge loss rate parameter α_{sh}
 - momentum conservation equation for zonal-averaged $\langle nv_y \rangle$
 - advances $\langle \Phi \rangle$ from $\langle v_y \rangle = \langle \partial \Phi / \partial x \rangle$
- important features
 - poloidal **directionality** from electron drift v_{*e}
 - radial **blob propagation** from curvature
 - radially varying dissipation (**sheaths**)

radial asymmetries drive
momentum flux:
v shear Gürcan et al., PoP (2007)
B shear
dissipation shear

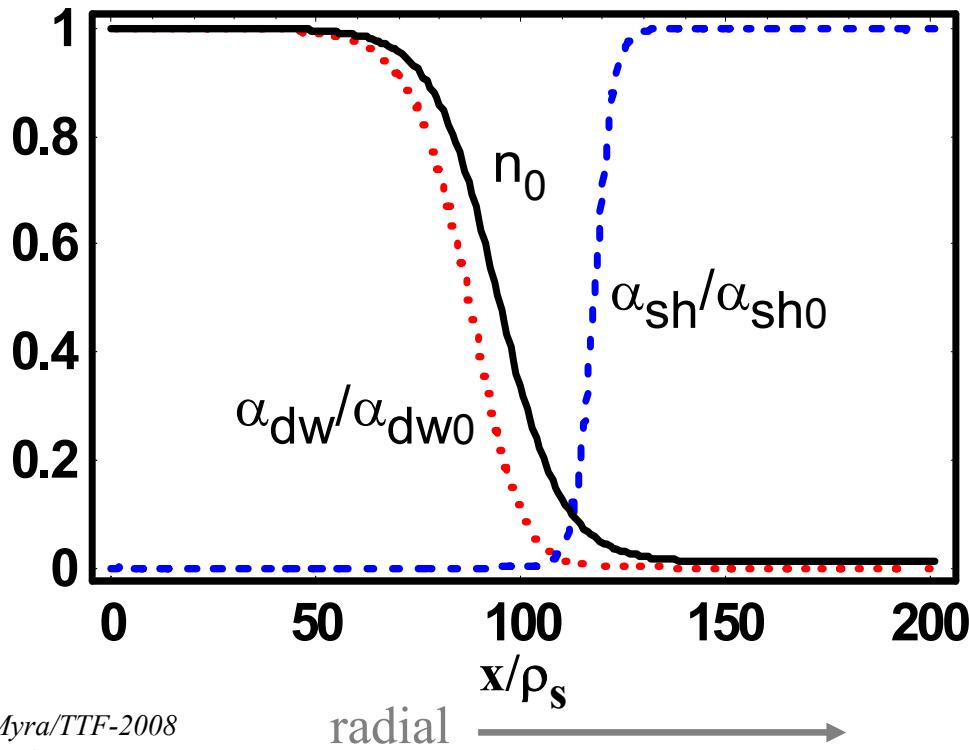
Equations and profiles

$$\frac{\partial n}{\partial t} + \mathbf{v} \cdot \nabla n = \alpha_{dw} (\tilde{\Phi} - \tilde{N}) - \alpha_{sh} n + D \nabla^2 n + S$$

$$\frac{\partial}{\partial t} \nabla^2 \tilde{\Phi} = \left\{ -\mathbf{v} \cdot \nabla \nabla^2 \Phi + \alpha_{dw} \left(\frac{\tilde{\Phi} - \tilde{N}}{n} \right) + \alpha_{sh} \tilde{\Phi} - \beta \frac{\partial \tilde{N}}{\partial y} + \mu \nabla^4 \tilde{\Phi} \right\}$$

$$\frac{\partial}{\partial t} \langle nv_y \rangle + \frac{\partial}{\partial x} \langle nv_x v_y \rangle = \int_0^x dx' \alpha_{sh} \langle n \Phi \rangle + \bar{\mu} \frac{\partial^2}{\partial x^2} \langle v_y \rangle$$

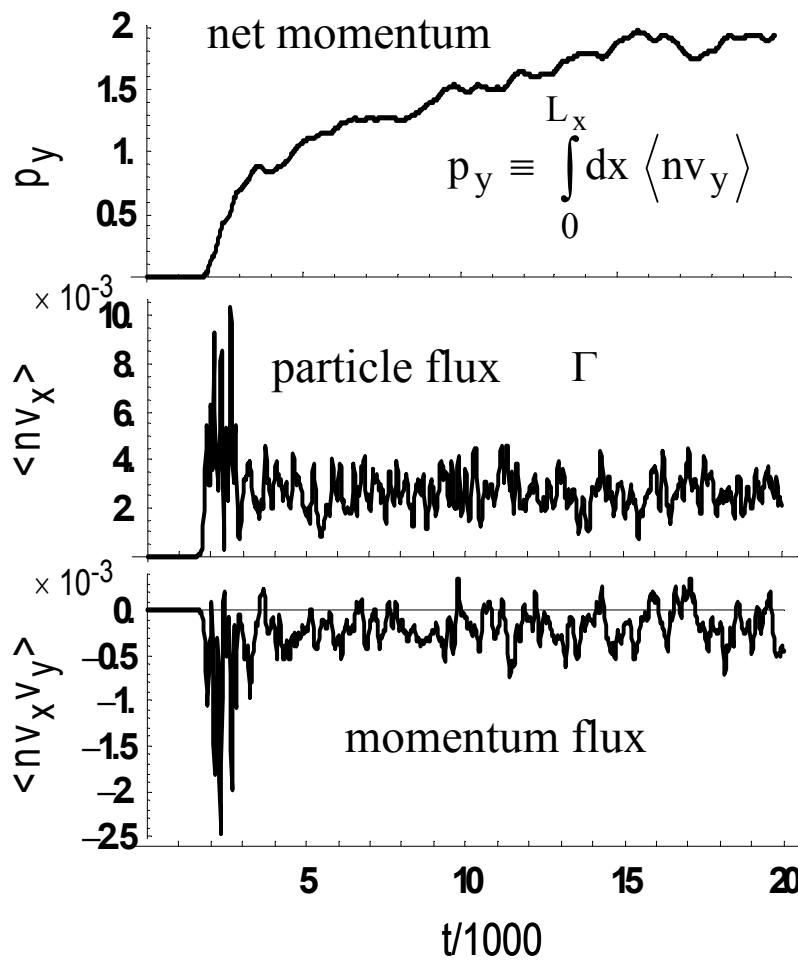
→ $\langle J_x \rangle = - \int_0^x dx' \langle \nabla_{||} J_{||} \rangle$



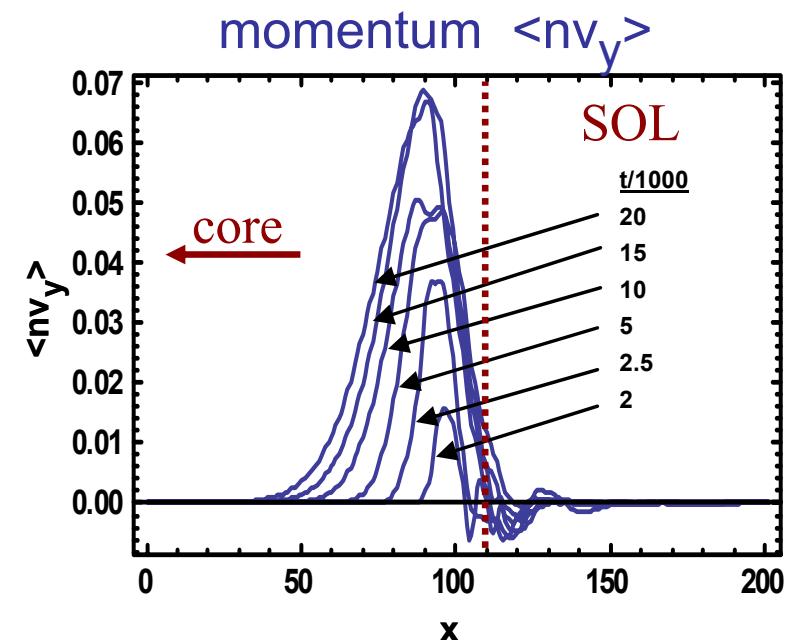
- closed surfaces (“edge”) have drift physics α_{dw}
- open SOL field lines have sheath physics α_{sh}
- curvature parameter $\beta =$ constant
- diffusion: $D, \mu, \bar{\mu}$

SOLT Code Results

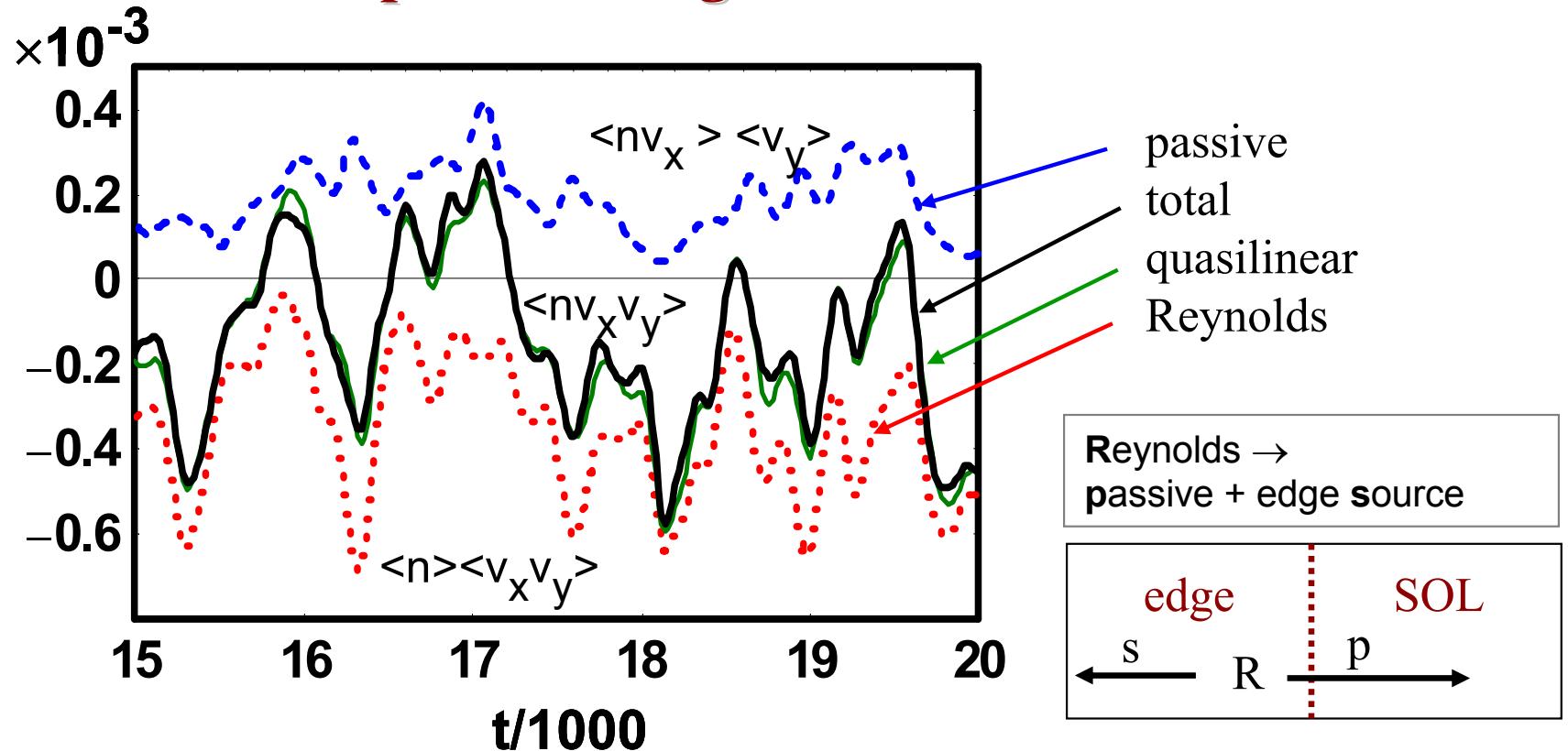
Ejection of moment across LCS and plasma spin-up



- net plasma momentum buildup (balanced by sheath momentum loss)
- initial transient blob “kick”
- intermittency
- inward momentum diffusion $\bar{\mu}$



Residual from Reynolds and passive momentum flux provides edge source



- full momentum flux \approx quasilinear average $\langle nv_x v_y \rangle \approx \langle v_y \rangle \langle nv_x \rangle + \langle n \rangle \langle v_x v_y \rangle$
- Reynolds and passive oppose
- Reynolds term generates flows, $\langle v_y \rangle$ builds up
 - passive momentum losses increase: v_y carried with the particle flux (blobs)
 - shear stabilization increases: v_y' vs. γ [see Russell et al., Edge IV]
- total, in direction of Reynolds (after cancellation) = edge source

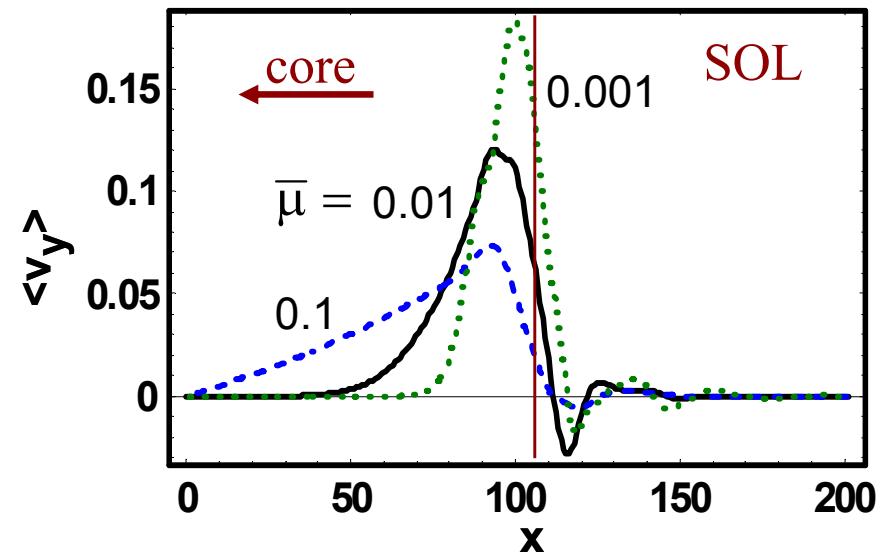
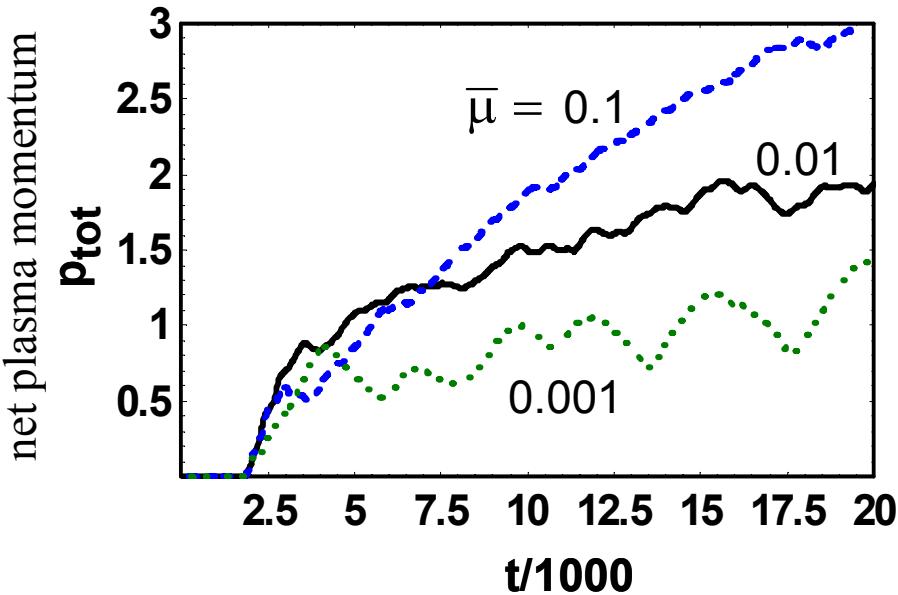
Edge momentum source strength depends on inward transport rate and edge velocity shear

- edge momentum transport into “core” is controlled by $\bar{\mu}$ in our model

$$\frac{\partial}{\partial t} \langle nv_y \rangle = \frac{\partial}{\partial x} \underbrace{\bar{\mu}}_{\text{inward momentum flux}} \frac{\partial \langle v_y \rangle}{\partial x} + \dots$$

~ NC or turb. pinch, diffn,
off-diag.

inward momentum flux



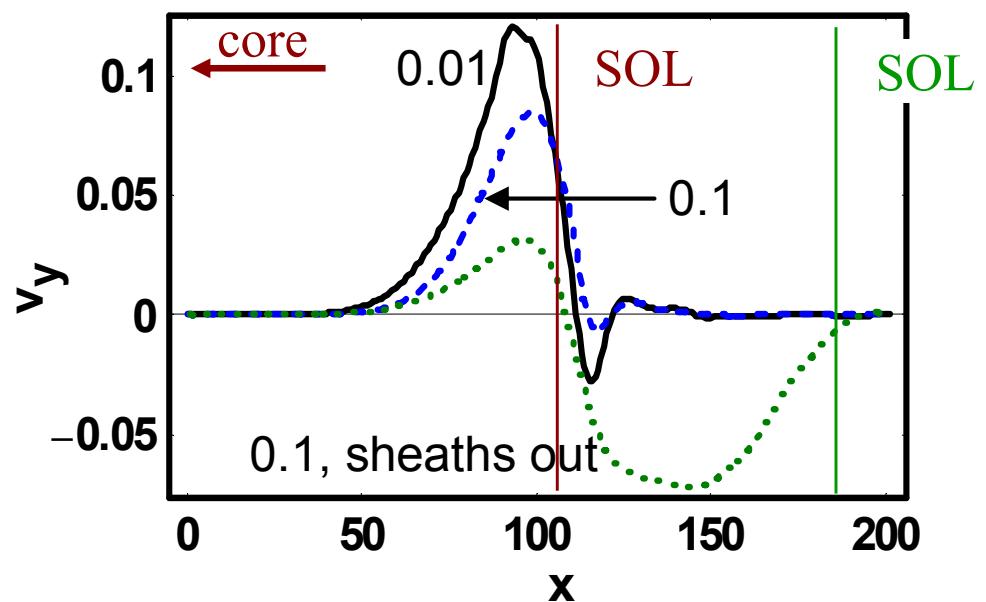
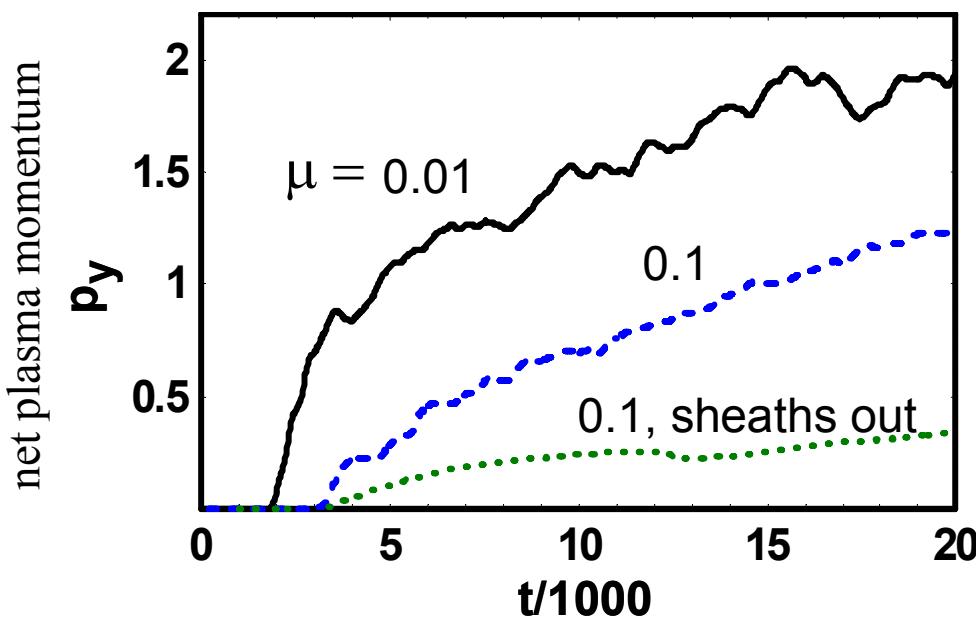
- larger inward transport \Rightarrow stronger edge source (to a point)
- system tries to maintain sheared flow in instability zone: $v_y' \sim \gamma$

nv_y vs. v_y'

edge – core coupling

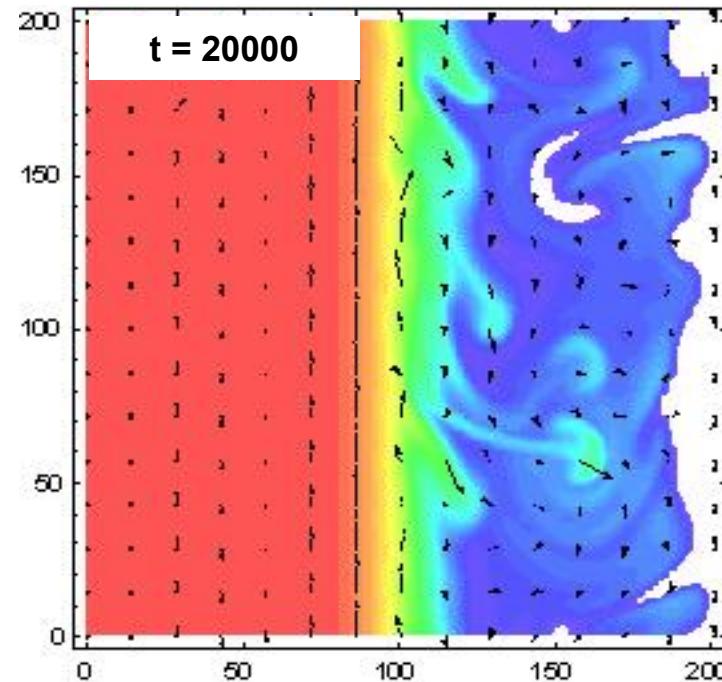
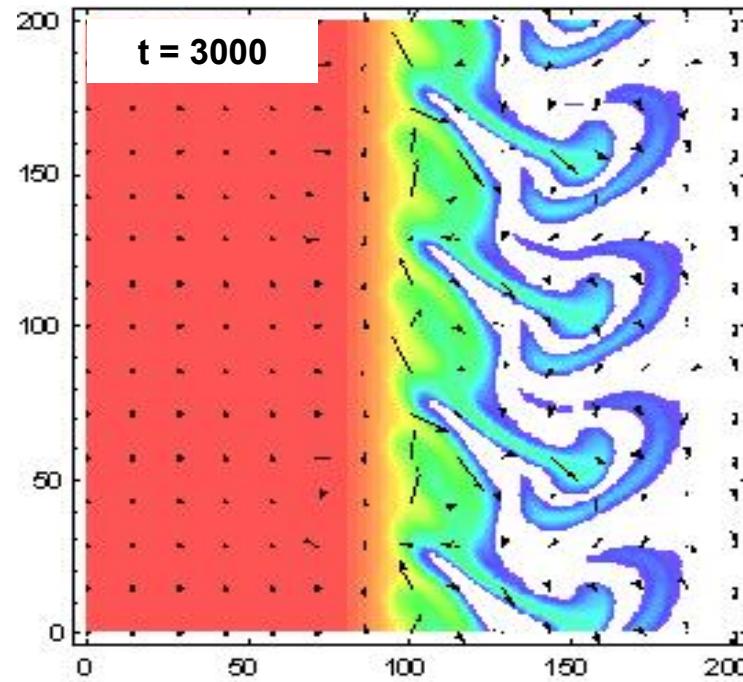
Viscous and sheath dissipation control the size and structure of the flows

- increased viscous dissipation (μ) of fluctuations reduces γ , turbulence and momentum transport
- artificially moving sheaths out (reduced sheath dissipation) \Rightarrow bipolar flows (less net plasma momentum)



Dynamics of blob formation and momentum transport

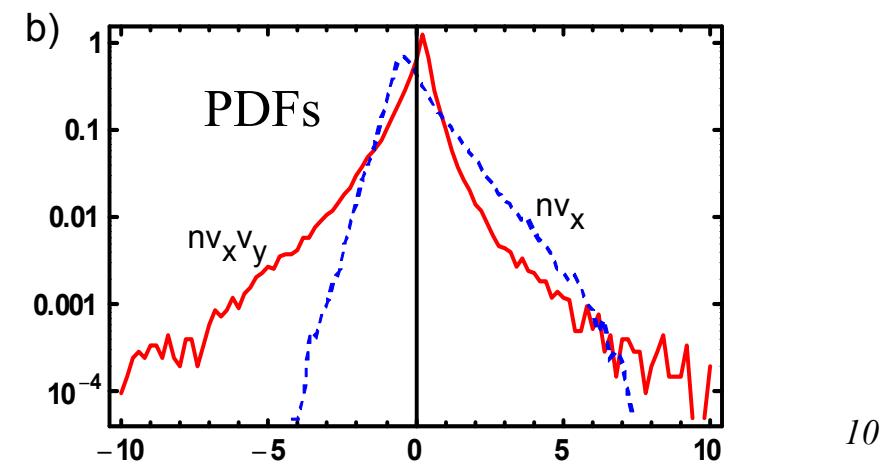
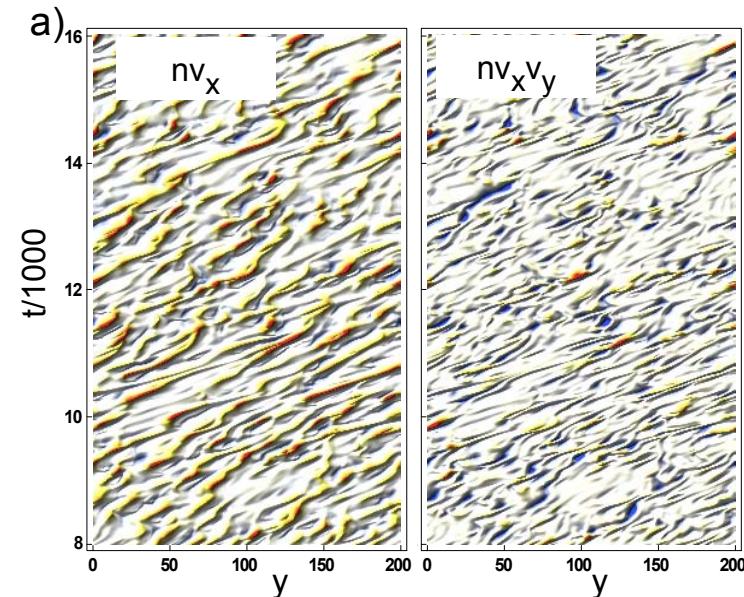
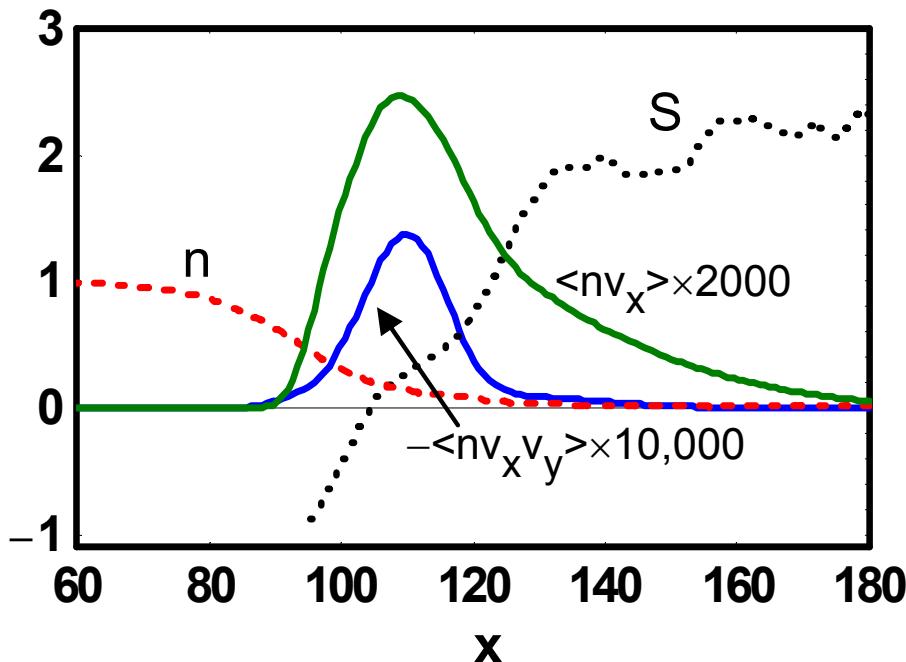
snapshots: density palette, momentum arrows



- early nonlinear development of seeded $m = 4$ mode
 - downward ejection of blobs (streamers); upward momentum “kick”
 - upward moving wave crests twisted around and down in ejection process
- later quasi-steady intermittency
 - sheared flows pinch off streamers \Rightarrow blobs

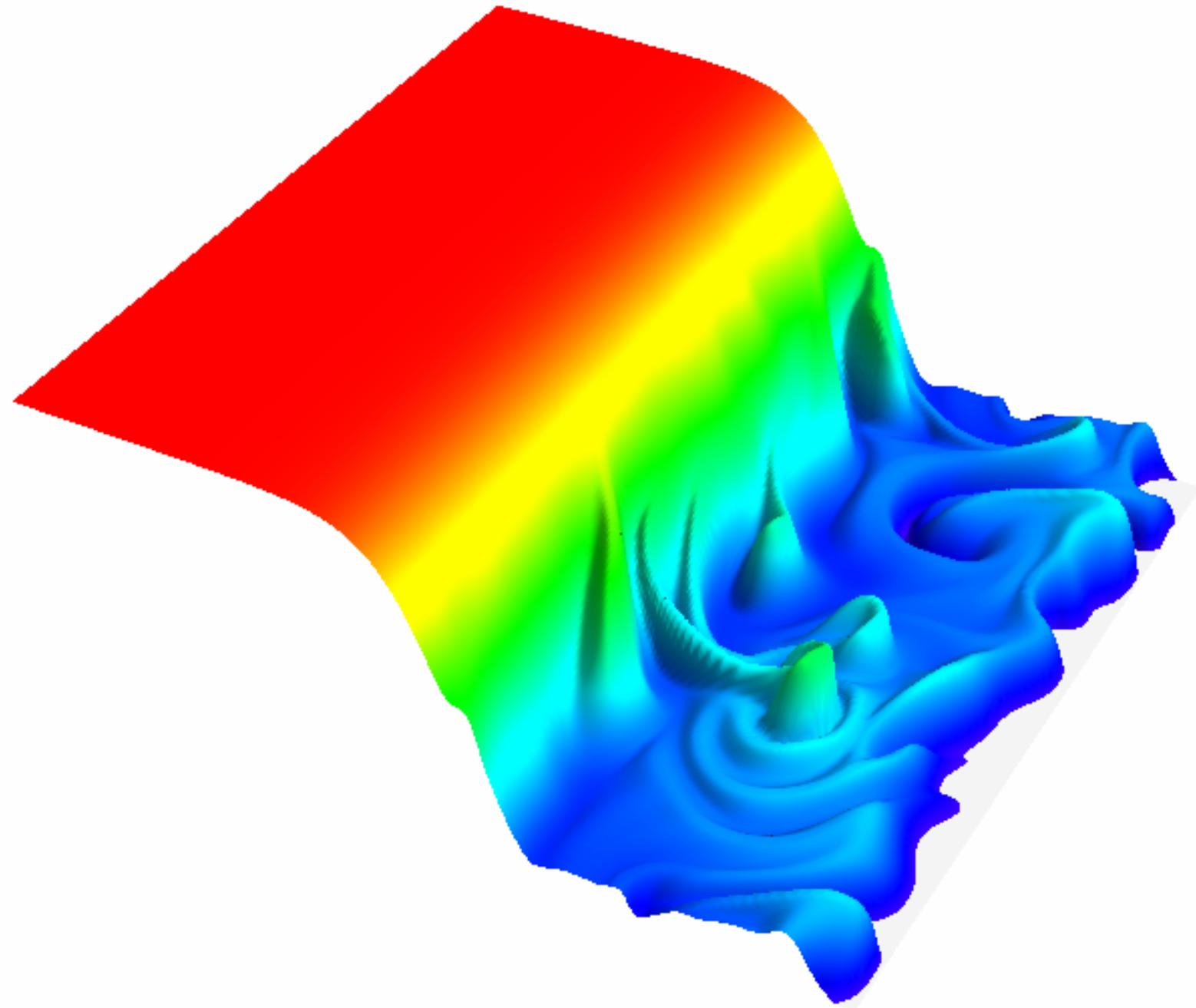
Density & momentum intermittency statistics

- momentum kick is mostly from blob birth zone where $|\text{skewness}| < 1$
- particle transport across LCS is by larger more coherent structures than momentum transport



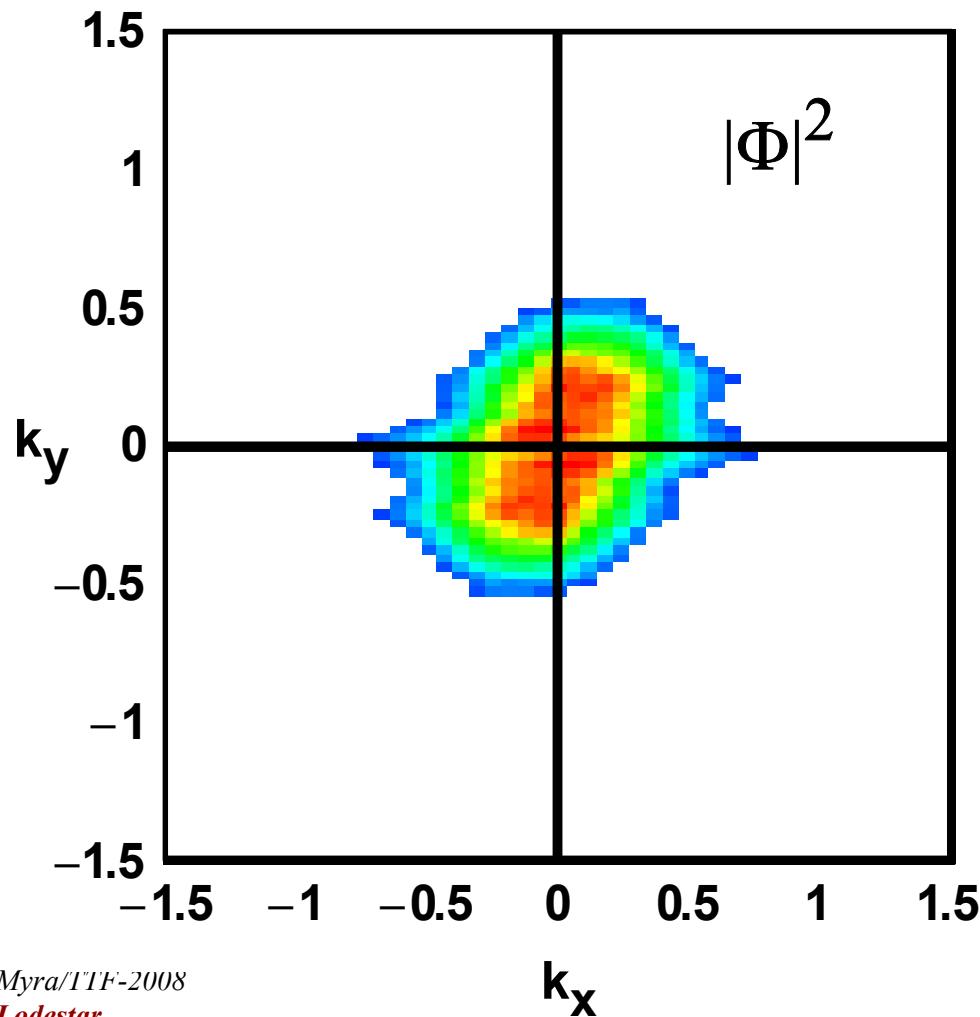
Conclusions & remarks

- *Poloidal* momentum transport studied here; *toroidal* momentum has some significant differences.
- Separatrix-spanning turbulence \Rightarrow edge momentum source for core rotation.
 - sheath momentum loss \Rightarrow net flow (vs. bipolar)
- Reynolds stress provides net force on the core plasma, and is opposed by the passive loss of momentum carried by exiting particles.
- Edge momentum source is coupled to core physics, inward transport rate.
- Steep edge $n(x)$ impacts dynamics:
 - regulation by edge sheared flows, viz. v_y' (no density weighting)
 - momentum conservation applied to nv_y
- Quasilinear averages work reasonably well; coherent structures are more important for particle transport than momentum.
- Turbulence-induced poloidal spin-up direction = direction of drift wave
 $v_{\text{phase}} \approx v_{*e}$



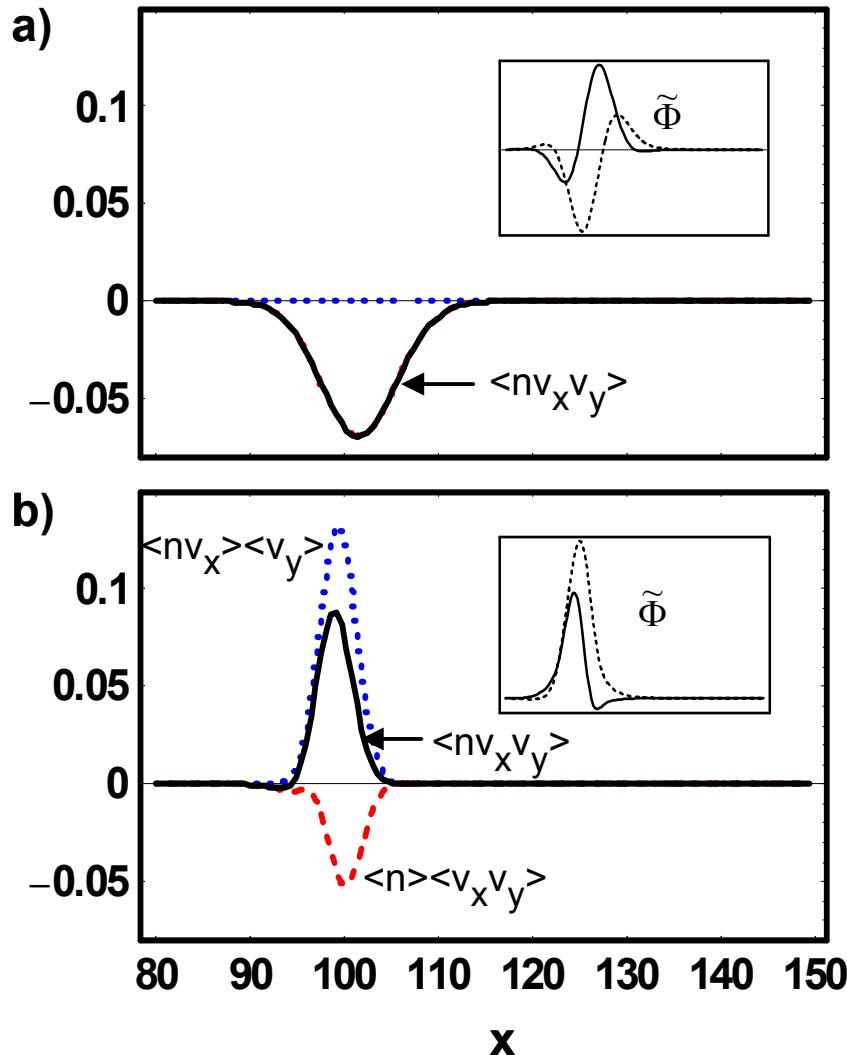
Supplemental

Spectral energy shows asymmetry necessary for momentum transport



- tilted ellipse necessary to get non-zero Reynolds stress
 $\langle v_x v_y \rangle \propto -k_x k_y \langle |\Phi|^2 \rangle$

Quasilinear averages provide a reasonable description



- total and constituent pieces of momentum flux from radial eigenvalue code + QL averages
- a) $v_y = 0$ (no passive term)
- b) large v_y (passive term dominates)
- critical v_y depends on ratio particle/momentum flux
 - QL does reasonable job

Toroidal (\approx parallel) momentum transport

- consider $\langle nv_x v_\zeta \rangle = \langle n \rangle \langle v_x v_\zeta \rangle + \langle nv_x \rangle \langle v_\zeta \rangle$

Reynolds passive



$$\langle \tilde{v}_x \tilde{v}_\zeta \rangle = -b_0 \langle \tilde{v}_x \tilde{v}_y \rangle + b_\zeta \langle \tilde{v}_x \tilde{v}_{||} \rangle$$

- need closure model for $\tilde{v}_{||}$

- v_{||} response of plasma to k_{||}

$$\tilde{v}_{||} = \frac{k_{||}}{\omega} c_s^2 \frac{\tilde{n}}{n} + \text{term } \propto \tilde{v}_x \partial v_{||} / \partial x$$

- resulting Reynolds term has diffusion + extra “pinch” effect

$$\langle \tilde{v}_x \tilde{v}_{||} \rangle = \frac{c_s^2}{nv_{||\text{ph}}} \langle \tilde{v}_x \tilde{n} \rangle + \dots$$



wave momentum carried by particles, blobs

- Coppi, Nucl. Fusion 42, 1 (2002).
- Coppi, IAEA Lyons, (2002).
- Gurcan et al., PoP (2007)
- Coppi et al, EPS 2006 paper P4.017**
- Myra et al., IAEA Chengdu (2006).